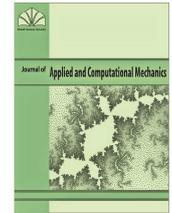




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Research Paper

Two Modifications of the Homotopy Perturbation Method for Nonlinear Oscillators

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Abstract: Nonlinear vibration arises in engineering and physics, and the periodic motion of these nonlinear oscillatory systems have rich dynamics. An estimation of amplitude-frequency relationship of a nonlinear oscillator is much needed, therefore, well-known homotopy perturbation method is employed for this purpose. In this paper, two last modifications of the homotopy perturbation method are briefly reviewed, which couples with either the parameter-expansion technology or the enhanced perturbation method. Both modifications are extremely effective for nonlinear oscillators, and the cubic-quintic-septic Duffing oscillator is used as an example to elucidate the solution processes.

Keywords: Homotopy perturbation method, Enhanced perturbation method, Amplitude-frequency relationship, Duffing oscillators, Periodic solution.

1. Introduction

Nonlinear oscillation from the oscillation of a molecule's vibration to the earthquake happens everywhere, and the periodic property performs a critical role in various oscillatory problems of science and engineering [1-7], for example, the attachment oscillator for the controlled manufacturing of nanofiber membranes [8], Fangzhu's oscillator for collection of water from air [9], harmonic oscillator for micro/nano structures [10], and release oscillator for delivery of ions [11]. In recent years, nonlinear oscillators have received extensive attention from scientific researchers especially after the invention of microelectromechanical systems [12-16]. The exact solutions of these vibratory systems play a crucial role in researching the properties and behavior of periodic motion of the aforesaid systems but identifying such exact solutions is very difficult. Although numerical solutions of such oscillatory systems with nonlinearities are easy to find, researchers are interested in finding the exact or at least analytic solutions to these problems because they have more detail which helps better insight into these systems. Therefore, amplitude-frequency formulation is an essential property of a nonlinear oscillator. There are many analytical techniques used in open literature for amplitude-frequency formulation, for example, the homotopy perturbation method [17-22], the variational iteration method [23-28], the Hamiltonian approach [29-30], He's frequency formulation [31-32] and the max-min approach [33].

The homotopy perturbation method was first introduced in 1998 [17], and it has been widely applied to solve various nonlinear problems accurately and efficiently [19-22, 34-38]. The method became famous after 2006 when the review article "some asymptotic methods for strongly nonlinear equations" was published in International Journal of Modern Physics B, which has been cited 1,345 times according to Clarivate's Web of Science (December 7, 2018), the review article gave a lucid as well as elementary introduction to the homotopy perturbation method [21]. On February 2008, Prof. Ganji told Science Watch (<http://archive.sciencewatch.com/dr/fbp/2008/08febfbp/08febGanji/>) that the homotopy perturbation method is a universal mathematical tool for solving nonlinear problems, and it will become as popular as Newton's iteration method in future. After 2006, publications increased exponentially and now there are about 200 publications per year, see Fig.1. Fractional differential equations [39-41] can be effectively solved by the homotopy perturbation method coupled with Laplace transform, this modification is called as He-Laplace method [42,43].

Though the homotopy perturbation method has been widely adopted, there are still space for further improvement. In this paper we briefly review two modifications using either the parameter expansion technology [44] or the enhanced perturbation method [45-46].



Total Publications

2,768 Analyze

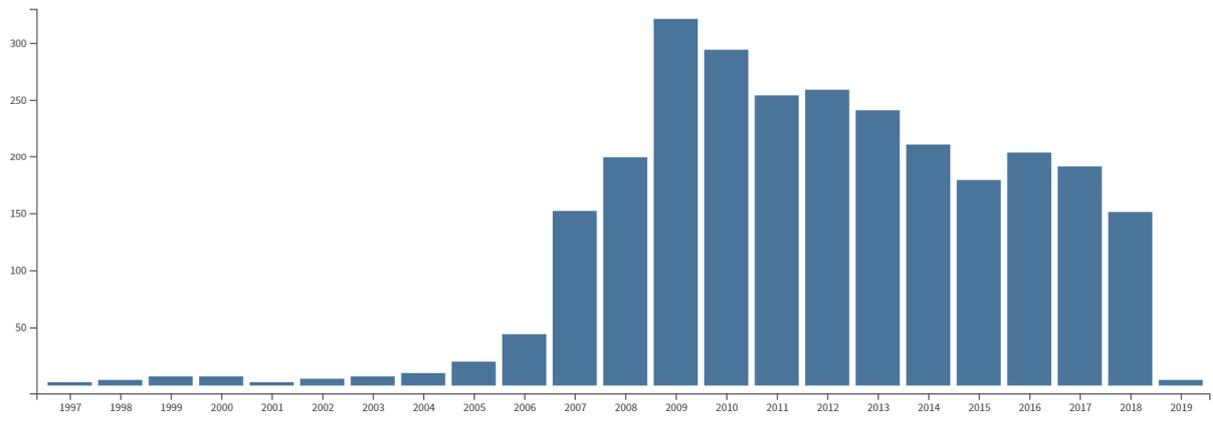


Fig. 1. Publication numbers on “homotopy perturbation method” per year according to Clarivate’s Web of Science (December 7, 2018)

2. The homotopy perturbation method with the parameter-expanding method

Consider the general non-linear oscillator equation as

$$x'' + ax + f(x', x) = 0 \quad (1)$$

where $f(x', x)$ is supposed to be a function with non-linear terms, a is a real number, it can be zero or negative. We can construct the homotopy equation for Eq.(1) as

$$x'' + ax + p[f(x', x)] = 0 \quad (2)$$

The solution can be written as

$$x = x_0 + px_1 + p^2x_2 + \dots \quad (3)$$

According to the parameter-expanding method [20,44] we can expand the coefficient of the linear term in the form

$$a = \Omega^2 + a_1p + a_2p^2 + \dots \quad (4)$$

where constants Ω^2 and a_i can be identified later.

3. Homotopy perturbation method coupled with the enhanced perturbation method

The perturbation method has been one of the powerful methods used to solve nonlinear problems. Roughly all perturbation theories build on the hypothesis of small parameters. Recently, Filobello-Nino et al. [45] suggested a new perturbation approach to solve the nonlinear equations in order to get the solution not influenced by parameters. Filobello-Nino et al. [45] proposed an adjustment in the perturbation method which was named as the enhanced perturbation method. This method is highly accurate and provides better results, and it can deal problems with both small and large values of the perturbation parameters [46]. To understand the idea of the enhanced perturbation method, we consider the following cubic quintic Duffing oscillator [47]

$$x'' + x + \varepsilon_1x^3 + \varepsilon_2x^5 = \cos \beta t \quad (5)$$

Eq. (7) can be expressed as

$$(D^2 + 1 + \varepsilon_1x^2 + \varepsilon_2x^4)x = \cos \beta t \quad (6)$$

where $D = d/dt$ is differential operator.

The enhanced perturbation method [31-32] is to apply the annihilator operator $D^2 + \beta^2$ to Eq. (6), we have

$$(D^2 + \beta^2)(D^2 + 1 + \varepsilon_1x^2 + \varepsilon_2x^4)x = 0 \quad (7)$$

This method can solve wide class of non-linear problems. It is more effective in case of non-linear problems with forced term but can also apply on problems without forced term [48].

4. Solution of cubic-quintic-septic duffing oscillator

The Duffing equation [47,49-50] is extensively studied non-linear differential equation and related to many engineering systems. In this article, this paper focuses on the cubic-quintic-septic Duffing equation [50]:

$$x'' + \alpha x + \beta x^3 + \gamma x^5 + \delta x^7 = 0 \quad (8)$$



with initial conditions

$$x(0) = A, \quad x'(0) = 0 \tag{9}$$

where α is coefficient of linear term and β, γ, δ are the cubic, quintic, septic power coefficients respectively. Eq. (8) expresses the Quasi-zero stiffness phenomena which is useful for designing the applicable structural parameters. This equation looks simple but difficult to solve analytically. The difficulty in obtaining accurate result of Eq. (8) is due to the presence of quintic and septic power non-linearities which make it more complex and complicated. Recently Remmi and Latha [50] obtained an approximate solution using the harmonic balancing method.

We can construct a homotopy equation for Eq. (8) as [20,44]

$$x'' + \alpha x + p[\beta x^3 + \gamma x^5 + \delta x^7] = 0 \tag{10}$$

The coefficient of linear term (α) and the solution can be expanded into following forms

$$\alpha = \Omega^2 + a_1 p + a_2 p^2 + \dots \tag{11}$$

$$x = x_0 + p x_1 + p^2 x_2 + \dots \tag{12}$$

Placing Eq. (11) and Eq.(12) into Eq. (10) and proceeding as that by the perturbation method, we have

$$x_0'' + \Omega^2 x_0 = 0 \quad x_0(0) = A, \quad x_0'(0) = 0 \tag{13}$$

$$x_1'' + \Omega^2 x_1 + x_0 a_1 + \beta x_0^3 + \gamma x_0^5 + \delta x_0^7 = 0, \quad x_1(0) = 0, \quad x_1'(0) = 0 \tag{14}$$

Solving x_0 from Eq. (13), we have

$$x_0 = A \cos \Omega t \tag{15}$$

Substituting Eq. (15) into Eq. (14), we obtain

$$\begin{aligned} &x_1'' + \Omega^2 x_1 + a_1 A \cos \Omega t + \beta A^3 \cos^3 \Omega t + \gamma A^5 \cos^5 \Omega t + \delta A^7 \cos^7 \Omega t = 0 \\ &x_1'' + \Omega^2 x_1 + \left(a_1 A + \frac{3}{4} \beta A^3 + \frac{5}{8} \gamma A^5 + \frac{35}{64} \delta A^7 \right) \cos \Omega t + \left(\frac{1}{4} \beta A^3 + \frac{5}{16} \gamma A^5 + \frac{21}{64} \delta A^7 \right) \cos 3\Omega t \\ &+ \left(\frac{1}{16} \gamma A^5 + \frac{7}{64} \delta A^7 \right) \cos 5\Omega t + \left(\frac{1}{64} \delta A^7 \right) \cos 7\Omega t = 0 \end{aligned} \tag{16}$$

No secular term in x_1 requires

$$a_1 A + \frac{3}{4} \beta A^3 + \frac{5}{8} \gamma A^5 + \frac{35}{64} \delta A^7 = 0 \tag{17}$$

If first order approximation is needed from Eq. (11) we have

$$a_1 = \alpha - \Omega^2 \tag{18}$$

Thus from Eq. (17) and Eq. (18), we have

$$\Omega = \sqrt{\alpha + \frac{3}{4} \beta A^2 + \frac{5}{8} \gamma A^4 + \frac{35}{64} \delta A^6} \tag{19}$$

In order to introduce the enhanced perturbation method [45-46] into the homotopy perturbation method, we first consider the following linear oscillator

$$x'' + \Omega^2 x = (D^2 + \Omega^2)x = 0 \tag{20}$$

According to the enhanced perturbation method [45-46], Eq. (20) can lead to a higher order form:

$$(D^2 + \Omega^2)(D^2 + \Omega^2)x = x'''' + 2\Omega^2 x'' + \Omega^4 x = 0 \tag{21}$$

Replacing x'' by $-\alpha \Omega^2 x$ from Eq. (20) results in

$$x'''' - \Omega^4 x = 0 \tag{22}$$

We adopt the similar way to increase the order of cubic-quintic-septic Duffing equation. To achieve this, differentiating Eq. (8) twice

$$x'''' + \alpha x' + 3\beta x^2 x' + 5\gamma x^4 x' + 7\delta x^6 x' = 0 \tag{23}$$

$$x'''' + \alpha x'' + 3\beta(x^2 x'' + 2x x'^2) + 5\gamma(x^4 x'' + 4x^3 x'^2) + 7\delta(x^6 x'' + 6x^5 x'^2) = 0 \tag{24}$$

Replacing x'' by $-(\alpha x + \beta x^3 + \gamma x^5 + \delta x^7)$ from Eq. (8), we have



$$x'''' - \alpha^2 x - \alpha\beta x^3 - \alpha\gamma x^5 - \alpha\delta x^7 + 3\beta(x^2 x'' + 2xx'^2) + 5\gamma(x^4 x'' + 4x^3 x'^2) + 7\delta(x^6 x'' + 6x^5 x'^2) = 0 \quad (25)$$

or

$$x'''' - \alpha^2 x - \alpha\beta x^3 - \alpha\gamma x^5 - \alpha\delta x^7 + 3\beta x^2 x'' + 6\beta x x'^2 + 5\gamma x^4 x'' + 20\gamma x^3 x'^2 + 7\delta x^6 x'' + 42\delta x^5 x'^2 = 0 \quad (26)$$

A homotopy equation can be established for Eq. (26) as

$$x'''' - \alpha^2 x + p[-\alpha\beta x^3 - \alpha\gamma x^5 - \alpha\delta x^7 + 3\beta x^2 x'' + 6\beta x x'^2 + 5\gamma x^4 x'' + 20\gamma x^3 x'^2 + 7\delta x^6 x'' + 42\delta x^5 x'^2] = 0 \quad (27)$$

The solution can be expanded as Eq. (12) and the coefficient of linear term can be expanded as

$$\alpha^2 = \Omega^4 + p\Omega_1 + p^2\Omega_2 + \dots \quad (28)$$

Substituting Eq. (12) and Eq. (28) into Eq. (27) and continuing as that by the perturbation method, we have

$$x_0'''' - \Omega^4 x_0 = 0, \quad x_0(0) = A, \quad x_0'(0) = 0 \quad (29)$$

$$x_1'''' - \Omega^4 x_1 - \Omega_1 x_0 - \alpha\beta x_0^3 - \alpha\gamma x_0^5 - \alpha\delta x_0^7 + 3\beta x_0^2 x_0'' + 6\beta x_0 x_0'^2 + 5\gamma x_0^4 x_0'' + 20\gamma x_0^3 x_0'^2 + 7\delta x_0^6 x_0'' + 42\delta x_0^5 x_0'^2 = 0, \quad x_1(0) = 0, \quad x_1'(0) = 0 \quad (30)$$

Using Eq. (15) as its initial approximate solution, we have

$$x_1'''' - \Omega^4 x_1 - A\Omega_1 \cos \Omega t - \alpha\beta A^3 \cos^3 \Omega t - \alpha\gamma A^5 \cos^5 \Omega t - \alpha\delta A^7 \cos^7 \Omega t - 3A^3 \Omega^2 \beta \cos^3 \Omega t + 6A^3 \Omega^2 \beta \cos \Omega t \sin^2 \Omega t - 5A^5 \Omega^2 \gamma \cos^5 \Omega t + 20A^5 \Omega^2 \gamma \cos^3 \Omega t \sin^2 \Omega t - 7A^7 \Omega^2 \delta \cos^7 \Omega t + 42A^7 \Omega^2 \delta \cos^5 \Omega t \sin^2 \Omega t = 0 \quad (31)$$

By means of simple calculation, we conclude

$$x_1'''' - \Omega^4 x_1 + (6A^3 \Omega^2 \beta - A\Omega_1) \cos \Omega t + (20A^5 \Omega^2 \gamma - 9A^3 \Omega^2 \beta - \alpha\beta A^3) \cos^3 \Omega t + (42A^7 \Omega^2 \delta - 25A^5 \Omega^2 \gamma - \alpha\gamma A^5) \cos^5 \Omega t - (49A^7 \Omega^2 \delta + \alpha\delta A^7) \cos^7 \Omega t = 0 \quad (32)$$

or after applying trigonometry formulas, we have

$$\begin{aligned} & x_1'''' - \Omega^4 x_1 + \left(-\frac{1}{64} \alpha\delta A^7 - \frac{49}{64} \delta A^7 \Omega^2 \right) \cos 7\Omega t \\ & + \left(-\frac{1}{16} \alpha\gamma A^5 - \frac{7}{64} \alpha\delta A^7 - \frac{25}{16} \gamma A^5 \Omega^2 - \frac{175}{64} \delta A^7 \Omega^2 \right) \cos 5\Omega t + \\ & \left(-\frac{1}{4} \alpha\beta A^3 - \frac{5}{16} \alpha\gamma A^5 - \frac{21}{64} \alpha\delta A^7 - \frac{9}{4} \beta A^3 \Omega^2 - \frac{45}{16} \gamma A^5 \Omega^2 - \frac{189}{64} \delta A^7 \Omega^2 \right) \cos 3\Omega t \\ & + \left(-A\Omega_1 - \frac{3}{4} \alpha\beta A^3 - \frac{5}{8} \alpha\gamma A^5 - \frac{35}{64} \alpha\delta A^7 - \frac{3}{4} \beta A^3 \Omega^2 - \frac{5}{8} \gamma A^5 \Omega^2 - \frac{35}{64} \delta A^7 \Omega^2 \right) \cos \Omega t = 0 \end{aligned} \quad (33)$$

Requirement of no secular term needs

$$-A\Omega_1 - \frac{3}{4} \alpha\beta A^3 - \frac{5}{8} \alpha\gamma A^5 - \frac{35}{64} \alpha\delta A^7 - \frac{3}{4} \beta A^3 \Omega^2 - \frac{5}{8} \gamma A^5 \Omega^2 - \frac{35}{64} \delta A^7 \Omega^2 = 0 \quad (34)$$

If the first order approximation is enough, then from Eq. (28), we have

$$\Omega_1 = \alpha^2 - \Omega^4 \quad (35)$$

Solving Ω from Eqs. (34) and (35), we obtain

$$\Omega = \sqrt{\alpha + \frac{3}{4} \beta A^2 + \frac{5}{8} \gamma A^4 + \frac{35}{64} \delta A^6} \quad (36)$$

which is same as that obtained by the homotopy perturbation method with the parameter- expanding method. Thus our approximate solution in both cases is

$$x(t) = A \cos \left(\sqrt{\alpha + \frac{3}{4} \beta A^2 + \frac{5}{8} \gamma A^4 + \frac{35}{64} \delta A^6} t \right) \quad (37)$$

Figure 2 shows the comparison of results obtained from both above explained methods with the numerical one with the help of MATLAB for the cubic-quintic-septic Duffing oscillator for some given parameters. We depict the approximate solution obtained from modifications of homotopy perturbation method (red circles) Eq. (37) with time for four different amplitudes $A = 0.05, 0.1, 0.15, 0.2$ against the parameter values $\alpha = \beta = \gamma = \delta = 1$ with the same yielded by computationally using MATLAB (solid black lines). All panels are showing good accuracy which validate the approximate solution obtained by homotopy perturbation method.

5. Conclusion

In this article, we successfully applied the two modifications of the homotopy perturbation method to the cubic-quintic-septic Duffing oscillator, both modifications result in a same result, giving a cross-proof of the correctness of the both modifications. Comparison with the numerical result show both modifications lead to a high accuracy of the obtained solution.



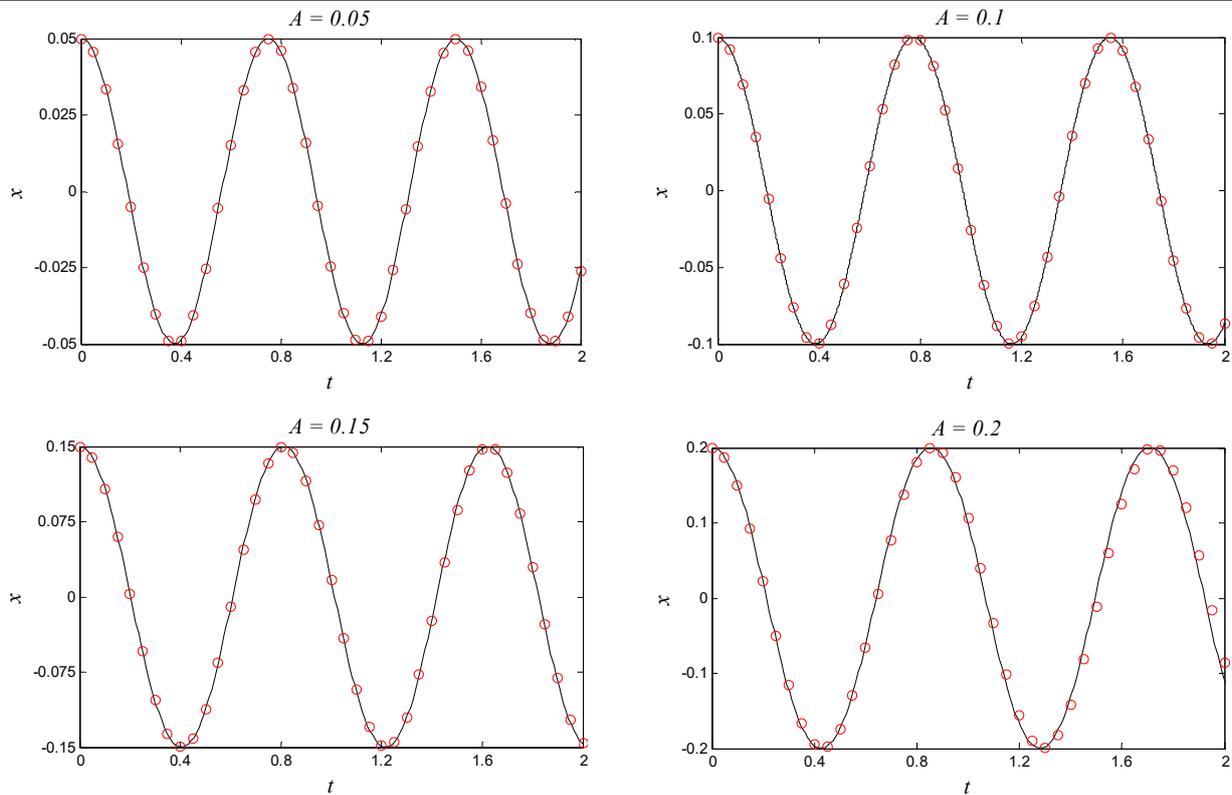


Fig. 2. Comparison of approximate solution with MATLAB solution for the cubic-quintic-septic oscillator

Author Contributions

All authors contribute equally in preparation of this manuscript. All authors discussed the results, reviewed and approved the final version of the manuscript.

Conflict of Interest

The authors declared no potential conflicts of interest with respect to the research, authorship, and publication of this article.

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Nomenclature

A	Amplitude of the oscillator	β	Cubic power coefficient
p	Expansion parameter	γ	Quintic power coefficient
x	Approximate solution	δ	Septic power coefficient
t	Time	Ω	Nonlinear frequency of the oscillator
α	Coefficient of linear term		

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